

Exercise 4.1

Displacement formulation of linear elastodynamics: strong and weak forms, Galerkin FE model

General IBVP of linear elastodynamics

Find $u_i = u_i(\mathbf{x}, t) = ?$, $\varepsilon_{ij} = \varepsilon_{ij}(\mathbf{x}, t) = ?$, $\sigma_{ij} = \sigma_{ij}(\mathbf{x}, t) = ?$ satisfying in Ω :

$$\sigma_{ij|j} + f_i = \rho \ddot{u}_i, \quad \varepsilon_{ij} = \frac{1}{2}(u_{i|j} + u_{j|i}), \quad \sigma_{ij} = \begin{cases} C_{ijkl} \varepsilon_{kl} & \text{(anisotropy),} \\ 2\mu \varepsilon_{ij} + \lambda \varepsilon_{kk} \delta_{ij} & \text{(isotropy),} \end{cases}$$

with the initial conditions (at $t = t_0$):

$$u_i(\mathbf{x}, t_0) = u_i^0(\mathbf{x}) \quad \text{and} \quad \dot{u}_i(\mathbf{x}, t_0) = v_i^0(\mathbf{x}) \quad \text{in } \Omega,$$

and subject to the boundary conditions on $\Gamma = \Gamma_u \cup \Gamma_t$ ($\Gamma_u \cap \Gamma_t = \emptyset$):

$$u_i(\mathbf{x}, t) = \hat{u}_i(\mathbf{x}, t) \quad \text{on } \Gamma_u, \quad \sigma_{ij}(\mathbf{x}, t) n_j = \hat{t}_i(\mathbf{x}, t) \quad \text{on } \Gamma_t.$$

1. Derive the displacement formulation of elasticity (Navier's equations)
2. Derive the weak variational form of the displacement formulation
3. Derive the Galerkin FE model (i.e., $\mathbf{M} \ddot{\mathbf{q}}(t) + \mathbf{K} \mathbf{q}(t) = \mathbf{Q}(t)$)

Exercise 4.2

A simple thermo-elastic problem

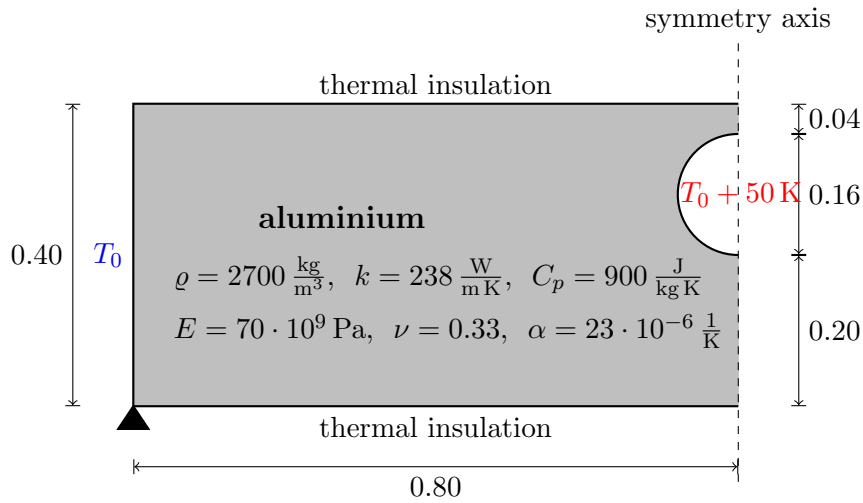


FIGURE 1: An uncoupled thermoelastic problem: solve the conductive heat transfer problem and use the temperature variation obtained in the domain as a thermal load for a coupled mechanical problem.

Exercise 4.3

Using the Weak Form PDE Interface in COMSOL Multiphysics

Weak form for harmonic linear elasticity

$$\omega^2 \int_{\Omega} \rho u_i \delta u_i - \int_{\Omega} \sigma_{ij} \delta u_{i|j} + \int_{\Omega} f_i \delta u_i + \int_{\Gamma_t} \hat{t}_i \delta u_i = 0$$

$$\text{Here: } \sigma_{ij} = \sigma_{ij}(\mathbf{u}) = \begin{cases} C_{ijkl} u_{k|l} & \text{– for anisotropic material} \\ \mu (u_{i|j} + u_{j|i}) + \lambda u_{k|k} \delta_{ij} & \text{– for isotropic material} \end{cases}$$

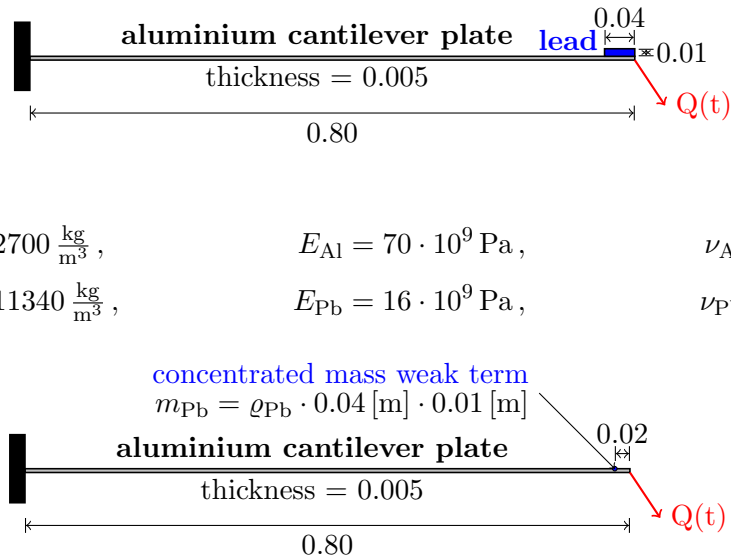


FIGURE 2: Simplified simulation of the influence of a heavy lead lump by adding a concentrated mass term to the weak form of a harmonic elastodynamic problem.

- Constants:

$$\begin{aligned} \omega &= 2\pi f \\ \mu &= E/2/(1+\nu) \quad \lambda = \nu E/(1+\nu)/(1-2\nu) \end{aligned}$$

- Domain expressions (for an isotropic material):

$$\begin{aligned} s_{11} &= 2\mu u_{1x} + \lambda(u_{1x} + u_{2y}) \\ s_{22} &= 2\mu u_{2y} + \lambda(u_{1x} + u_{2y}) \\ s_{12} &= \mu(u_{1y} + u_{2x}) \end{aligned}$$

- Domain integrand (in Ω):

$$\begin{aligned} &\omega^2 \rho (u_1 \text{test}(u_1) + u_2 \text{test}(u_2)) \\ &- s_{11} \text{test}(u_{1x}) - s_{12} \text{test}(u_{1y} + u_{2x}) - s_{22} \text{test}(u_{2y}) \\ &+ f_1 \text{test}(u_1) + f_2 \text{test}(u_2) \end{aligned}$$

- Neumann boundary integrand (on Γ_t):

$$t_1 \text{test}(u_1) + t_2 \text{test}(u_2)$$

- Concentrated mass weak term:

$$\omega^2 \text{mass} (u_1 \text{test}(u_1) + u_2 \text{test}(u_2))$$